

## A Molecular Orbital Study of the Conformation and $g$ -Factors of the HSSH<sup>-</sup> Radical Anion

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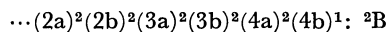
**Synopsis.** The electronic structure and  $g$ -factors of the HSSH<sup>-</sup> radical anion were calculated by means of a restricted open-shell SCF MO method with the MINDO/3 approximation, and the dependence of the  $g$ -factors upon the conformation was clarified. The most stable conformation predicted agreed well with the experimentally estimated -S-S- geometry of dialkyl disulfide anions, and its  $g$ -factors agreed well with the observed values.

Attention has been drawn in recent literature<sup>1,2)</sup> to the electronic structures of sulfur-containing radicals and their ESR data. Lin and Lunsford<sup>1)</sup> observed the HSSH<sup>-</sup> radical anion formed by allowing H<sub>2</sub>S to react with trapped electrons on MgO. From the hyperfine splittings and the  $g$ -factors ( $g_{\perp}=2.015$ ,  $g_{\parallel}=2.003$ ), they pointed out that the radical earlier identified as H<sub>2</sub>S<sup>-</sup> by Bennett *et al.*<sup>3)</sup> is HSSH<sup>-</sup>. Since the principal values of the  $g$ -tensor correlate strongly with the electronic structure of a radical, the theoretical analysis of the  $g$ -tensor can clarify the electronic structure and conformation of the radical. The purpose of the present note is to investigate the relation between the  $g$ -factors and the conformation of HSSH<sup>-</sup> by means of restricted open-shell SCF calculations.

The MO's of the spin-restricted doublet state were obtained by the approximate SCF version which has been proposed by Longuet-Higgins and Pople<sup>4)</sup> for  $\pi$ -radicals and which has since been applied to all-valence-electron systems by the present author.<sup>5)</sup> For the evaluation of the electron integrals involved, the MINDO/3 approximation<sup>6)</sup> was employed. The geometry of HSSH<sup>-</sup> was first optimized with respect to three parameters; the S-S bond length, the HSS angle,  $\phi$ , and the dihedral angle,  $\theta$ , between two planes containing the HSS' or SS'H group. The dependence of the  $g$ -factors upon  $\phi$  and  $\theta$  was calculated according to Stone's formula,<sup>7)</sup> the S-S bond length being fixed at the optimized value, 2.17 Å. A more detailed description of the calculation has been reported elsewhere.<sup>8)</sup>

### Results and Discussion

The electronic structure and  $g$ -factors of a typical conformation (shown in Fig. 1) will be discussed first, since the  $g$ -tensor of this conformation is practically diagonal in terms of the axes,  $x$ ,  $y$ , and  $z$ . The ground electron configuration of C<sub>2</sub> symmetry is:



where 2a and 2b MO's are bonding S-H  $\sigma$  orbitals and where 3a is the bonding S-S  $3p\sigma$  orbital directed along the S-S bond. The 3b and 4a orbitals are nearly degenerate and are expressed as  $(3p_{\alpha}(S)+3p_{\beta}(S'))/\sqrt{2}$  and  $(3p_{\alpha}(S)-3p_{\beta}(S'))/\sqrt{2}$  respectively. The half-occupied MO, 4b, is the anti-bonding  $3p\sigma^*$  orbital directed along the S-S bond. The contribution of

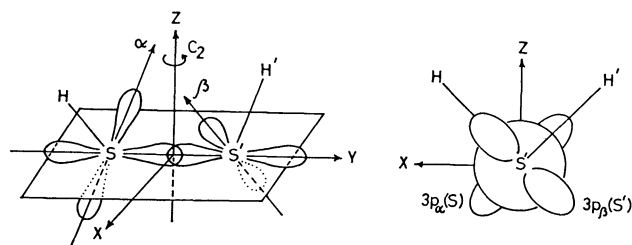


Fig. 1. The coordinate axes and the Newman projection for a typical conformation of HSSH<sup>-</sup> viewed in the direction of the S-S bond. The dihedral angle,  $\theta$ , between two planes containing H-S-S' and S-S'-H' is 90° and the HSS' angle,  $\phi$ , is 90°.

TABLE 1. CALCULATED  $\Delta g_{ii}^{(a)}$  VALUES OF THE HSSH<sup>-</sup> OF THE C<sub>2</sub> SYMMETRY, WITH  $\theta=90^\circ$  AND  $\phi=90^\circ$

Excited configuration	$\Delta E$ (eV) <sup>b)</sup>	$\Delta g_{xx}$	$\Delta g_{yy}$	$\Delta g_{zz}$
1a→4b	16.40	0	0	0
1b→4b	11.46	0	0	2
2a→4b	6.33	77	0	0
2b→4b	6.06	0	0	79
3a→4b	3.98	5	0	0
3b→4b	2.53	0	0	81
4a→4b	2.51	79	0	0
4b→5a	2.26	0	0	-66
4b→5b	2.29	-65	0	0
Total <sup>c)</sup>		96	0	96

a) Deviation from the free-spin value. The value of  $\Delta g_{ii} \times 10^4$  are listed in each column. b) Energy difference between the ground state and an excited configuration. c) The three principal values calculated from these values are:  $g_{xx}=2.0119$ ,  $g_{yy}=2.0023$ , and  $g_{zz}=2.0119$ .

each excited configuration to the principal values of the  $g$ -tensor is shown in Table 1. The large deviation of  $g_{xx}$  from the free-spin value can be understood as resulting from the spin-orbit mixing of the ground state with the 4a→4b excited configuration, and that of  $g_{zz}$  from the mixing with the 3b→4b excited configuration. The contributions of the bonding S-H  $\sigma$  orbitals (2a and 2b) compensate for those of the anti-bonding S-H  $\sigma^*$  orbitals (5a and 5b). The calculated principal values ( $g_{\perp}=2.011$  and  $g_{\parallel}=2.002$ ) are in fair agreement with the observed values.<sup>1)</sup>

The ground state of the HSSH<sup>-</sup> with the C<sub>2</sub> symmetry was calculated to be the <sup>2</sup>B state for any value of  $\phi$  or  $\theta$ . Though the half-occupied MO is always expressed as  $3p\sigma^*$ , its shape depends considerably upon  $\phi$  and  $\theta$ . This is the reason why the  $g$ -factors of HSSH<sup>-</sup> depend largely on the conformation of HSSH<sup>-</sup>. The dependence of  $g_{av}=1/3(g_1+g_2+g_3)$  upon  $\phi$  and  $\theta$  is shown in Fig. 2. It can be seen from Fig. 2 that there are many conformations for a given value of  $g_{av}$ . The three principal values, however, vary largely with the

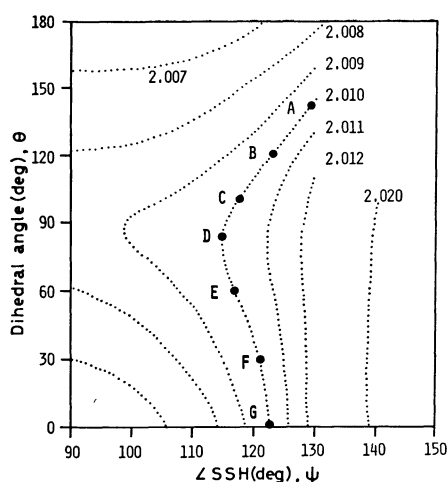


Fig. 2. The conformational dependence of the  $g_{av}$  value of HSSH $^-$ . The bond lengths shown in Fig. 4 were employed.

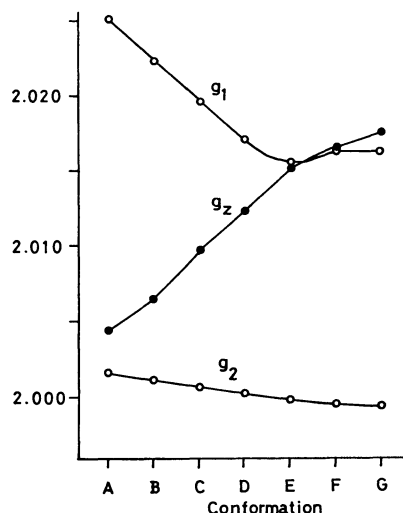


Fig. 3. Variations of three principal values of the  $g$ -tensor of HSSH $^-$  along A→G cited in Fig. 2.

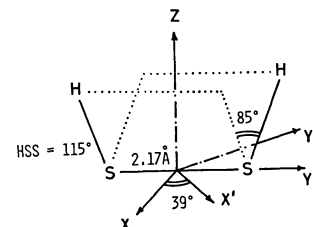


Fig. 4. The most stable conformation of HSSH $^-$  and the principal axes,  $x'$ ,  $y'$ , and  $z$ , of its  $g$ -tensor.

conformation of HSSH $^-$ . When  $g_{av}=2.010$ , for example the three principal values of the conformations corresponding to A—G in Fig. 2 vary as is shown in Fig. 3. From a detailed examination of the three principal values, it was found that the conformations whose calculated three principal values agree with the observed ones within  $\pm 0.003$  have (i)  $\theta=90^\circ$  and  $90^\circ < \psi < 115^\circ$  or (ii)  $60^\circ < \theta < 90^\circ$  and  $105^\circ < \psi < 115^\circ$ . The  $g$ -factors for the D conformation ( $\theta=85^\circ$ ,  $\psi=115^\circ$ ,  $SS=2.17$  Å,  $SH=1.358$  Å (ass.)) cited in Fig. 4 are worth listing, since this conformation is most stable in the present calculations and closely resembles the estimated  $-S-S-$  geometry of dialkyl disulfide anions.<sup>9)</sup> The  $g$ -factors calculated ( $g_{x'x'}=2.0172$ ,  $g_{y'y'}=2.0003$ ,  $g_{zz}=2.0124$ ,  $g_{av}=2.0010$ ) agree well with the observed values.<sup>1)</sup>

The radicals which have been attributed to  $R_2S_2^-$  are  $RCH_2-S-S-CH_2R^-$  ( $g_1=2.018$ ,  $g_{II}=2.002$ ),<sup>9)</sup> the  $n$ -amyl disulfide anion ( $g_1=2.022$ ,  $g_{II}=2.003$ ),<sup>10)</sup> and the dibenzoyldisulfide anion ( $g_1=2.014$ ,  $g_2=2.010$ ,  $g_3=2.004$ ).<sup>11)</sup> In spite of the large substituents in these radicals, their  $g$ -factors have the same features as those of HSSH $^-$ . This indicates that the dependence of the conformation of  $R_2S_2^-$  and of the electronic structure of the  $-S-S-$  group upon the substituent is relatively small, and that the present analysis of the  $g$ -factors gives information not only about HSSH $^-$ , but also about RSSR $^-$ .

The spin distribution of HSSH $^-$  is worth discussing since the identification of HSSH $^-$  has thus far rested on the  $^1H$  and  $^{33}S$  hyperfine coupling constants. The small  $^1H$  coupling constants<sup>1)</sup> suggest that the primary interaction in  $^1H$  coupling is due to spin polarization by an unpaired electron on the sulfur atom. The spin densities were thus calculated by means of the unrestricted SCF version<sup>12)</sup> with the MINDO/3 approximation. The spin density on the  $3s(S)$  AO is almost constant (0.06—0.07) for all the above conformations (i) and (ii). The spin density on the  $1s(H)$  AO depends largely on the HSS angle ( $-0.050(90^\circ)$ ,  $-0.040(100^\circ)$ ,  $-0.009(115^\circ)$ ,  $+0.001(120^\circ)$ , with the constant value of

$\theta=90^\circ$ ), while its dependence on  $\theta$  is small. The calculated spin distribution<sup>13)</sup> and observed  $^1H$  coupling constants (6—9 G)<sup>1)</sup> predict that the HSS angle will be  $105^\circ$ — $115^\circ$ .

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